

26th of June 2014 — Lattice 2014 (NYC, USA)

Inclusion of isospin breaking effects in lattice simulations

Antonin J. Portelli
(University of Southampton)

What's new ?

What's new ?

- ❖ [MILC, 2014] — **C. Bernard parallel talk: tomorrow 14:35**
 - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
 - new insights on finite-volume effects

What's new ?

- ❖ [MILC, 2014] — **C. Bernard parallel talk: tomorrow 14:35**
 - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
 - new insights on finite-volume effects
- ❖ [QCDSF, 2014] (pure QCD) — **R. Horsley parallel talk: today 14:55**
 - study of the $\Sigma^0 - \Lambda^0$ system

What's new ?

- ❖ [MILC, 2014] — **C. Bernard parallel talk: tomorrow 14:35**
 - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
 - new insights on finite-volume effects
- ❖ [QCDSF, 2014] (pure QCD) — **R. Horsley parallel talk: today 14:55**
 - study of the $\Sigma^0 - \Lambda^0$ system
- ❖ [BMWc, 2014] (EQ)
 - update of quark masses and Dashen's theorem using electro-quenched simulations

What's new ?

- ❖ [MILC, 2014] — **C. Bernard parallel talk: tomorrow 14:35**
 - update of quark masses and Dashen's theorem corrections using electro-quenched simulations
 - new insights on finite-volume effects
- ❖ [QCDSF, 2014] (pure QCD) — **R. Horsley parallel talk: today 14:55**
 - study of the $\Sigma^0 - \Lambda^0$ system
- ❖ [BMWc, 2014] (EQ)
 - update of quark masses and Dashen's theorem using electro-quenched simulations
- ❖ [Davoudi & Savage, 2014] — [arXiv:1402.6741]
 - finite-volume corrections to hadron masses in NREFTs

What's new ?

What's new ?

- ❖ [QCDSF, 2014] — **G. Schierholz parallel talk: tomorrow 14:15**
 - new full $N_f = 1+1+1$ QCD+QED simulations
 - preliminary results for the baryon octet splittings

What's new ?

- ❖ [QCDSF, 2014] — **G. Schierholz parallel talk: tomorrow 14:15**
 - new full $N_f = 1+1+1$ QCD+QED simulations
 - preliminary results for the baryon octet splittings
- ❖ [BMWc, 2014] — [arXiv: 1406.4088]
 - new set of $N_f = 1+1+1+1$ full QCD+QED simulations
 - extensive analytical/numerical study of finite-volume effects
 - high precision computation of the hadron spectrum splittings (continuum, infinite volume and physical point extrapolation)

- ❖ Motivations
- ❖ Update on electro-quenched results
- ❖ Lattice QED
- ❖ Full QCD+QED simulations
- ❖ Isospin splittings in the hadron spectrum
- ❖ Summary & outlook

Motivations

Isospin symmetry breaking

- ❖ Isospin symmetric world: up and down quarks are particles with identical physical properties.

Isospin symmetry breaking

- ❖ Isospin symmetric world: up and down quarks are particles with identical physical properties.

	up	down
Mass (MeV)	2.3 $\left(\begin{smallmatrix} +0.7 \\ -0.5 \end{smallmatrix}\right)$	4.8 $\left(\begin{smallmatrix} +0.5 \\ -0.3 \end{smallmatrix}\right)$
Charge (e)	2/3	-1/3

source: [PDG, 2013]

Isospin symmetry breaking

❖ Isospin symmetric world: up and down quarks are particles with identical physical properties.

❖ Isospin symmetry is explicitly broken by:

- the up and down **quark mass difference**

$$|m_u - m_d|/\Lambda_{\text{QCD}} \simeq 0.01$$

- the up and down **electric charge difference**

$$\alpha \simeq 0.0073$$

	up	down
Mass (MeV)	2.3 ^(+0.7) _(-0.5)	4.8 ^(+0.5) _(-0.3)
Charge (e)	2/3	-1/3

source: [PDG, 2013]

Nucleon mass splitting

❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

Nucleon mass splitting

- ❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

- ❖ needed for **proton stability**

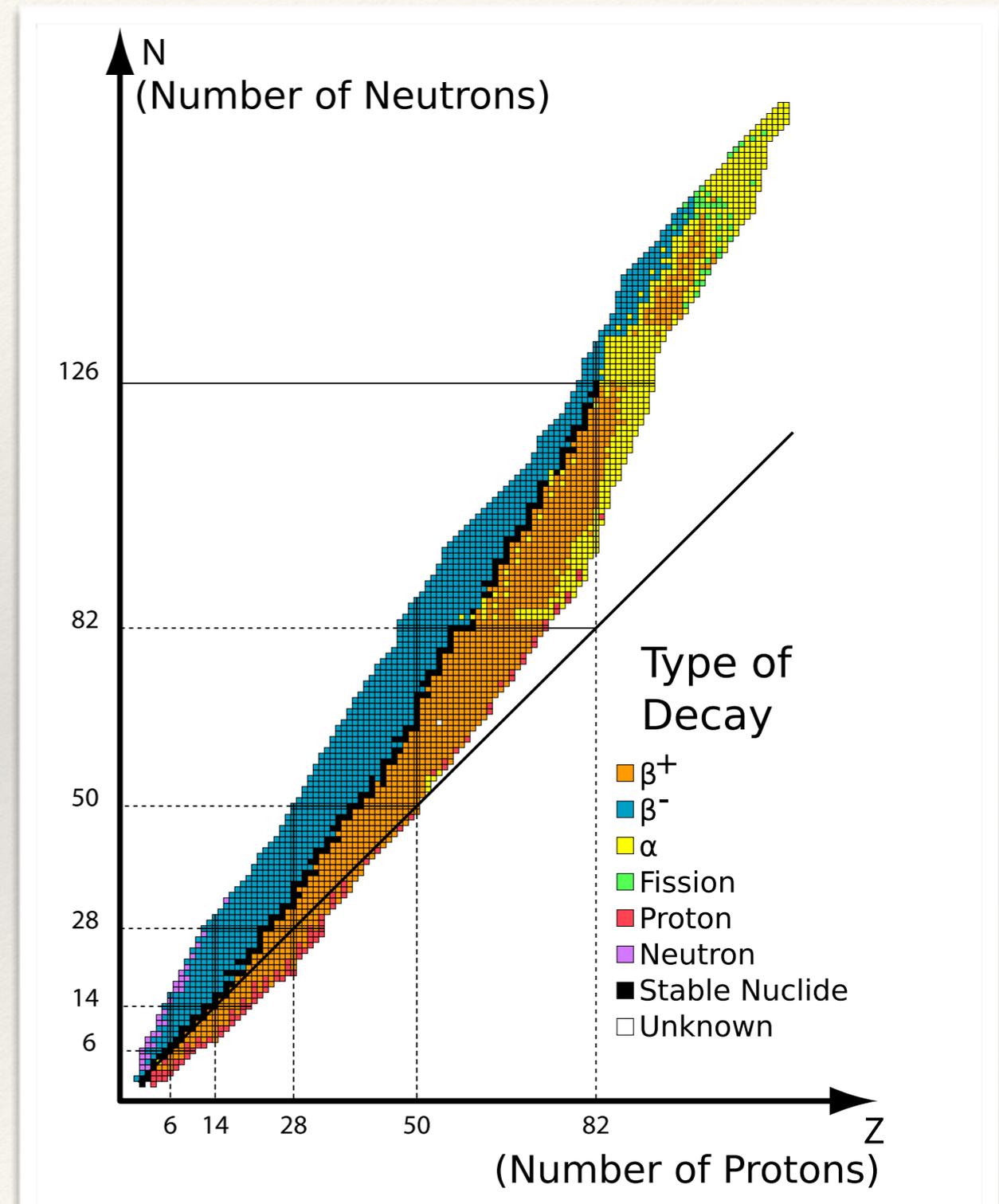
Nucleon mass splitting

- ❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

- ❖ needed for **proton stability**
- ❖ determines through β -decay the **stable nuclide chart**



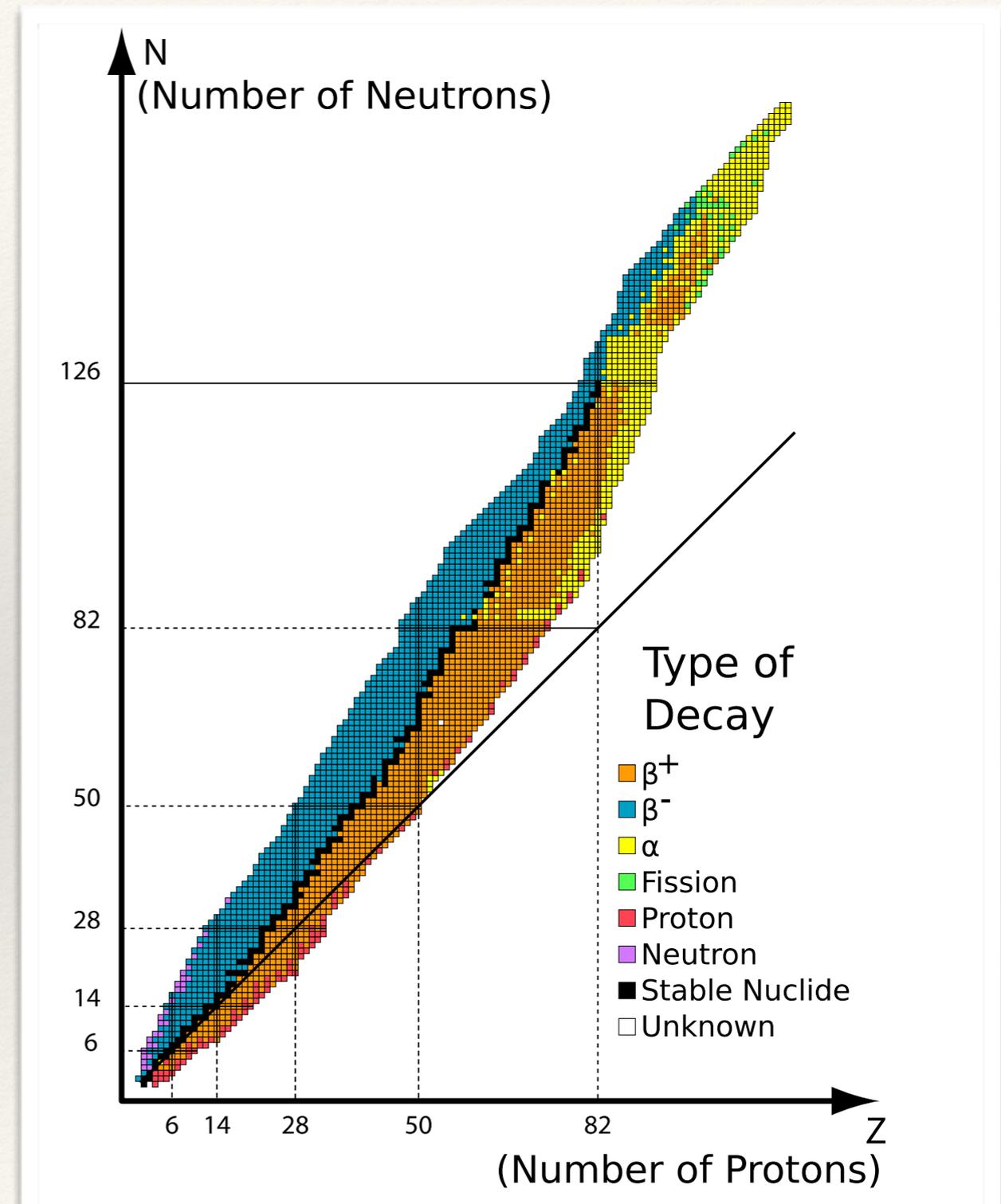
Nucleon mass splitting

- ❖ Well known experimentally:

$$M_n - M_p = 1.2933322(4) \text{ MeV}$$

source: [PDG, 2013]

- ❖ needed for **proton stability**
- ❖ determines through β -decay the **stable nuclide chart**
- ❖ initial condition for **Big-Bang nucleosynthesis**



Dashen's theorem

- ❖ In the SU(3) chiral limit [Dashen, 1969]:

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2 + \mathcal{O}(\alpha m_s)$$

Dashen's theorem

- ❖ In the SU(3) chiral limit [Dashen, 1969]:

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2 + \mathcal{O}(\alpha m_s)$$

- ❖ How large are the corrections? FLAG parametrisation:

$$\varepsilon = \frac{\Delta_{\text{QED}} M_K^2 - \Delta_{\text{QED}} M_\pi^2}{\Delta M_\pi^2}$$

Dashen's theorem

- ❖ In the SU(3) chiral limit [Dashen, 1969]:

$$\Delta_{\text{QED}} M_K^2 = \Delta_{\text{QED}} M_\pi^2 + \mathcal{O}(\alpha m_s)$$

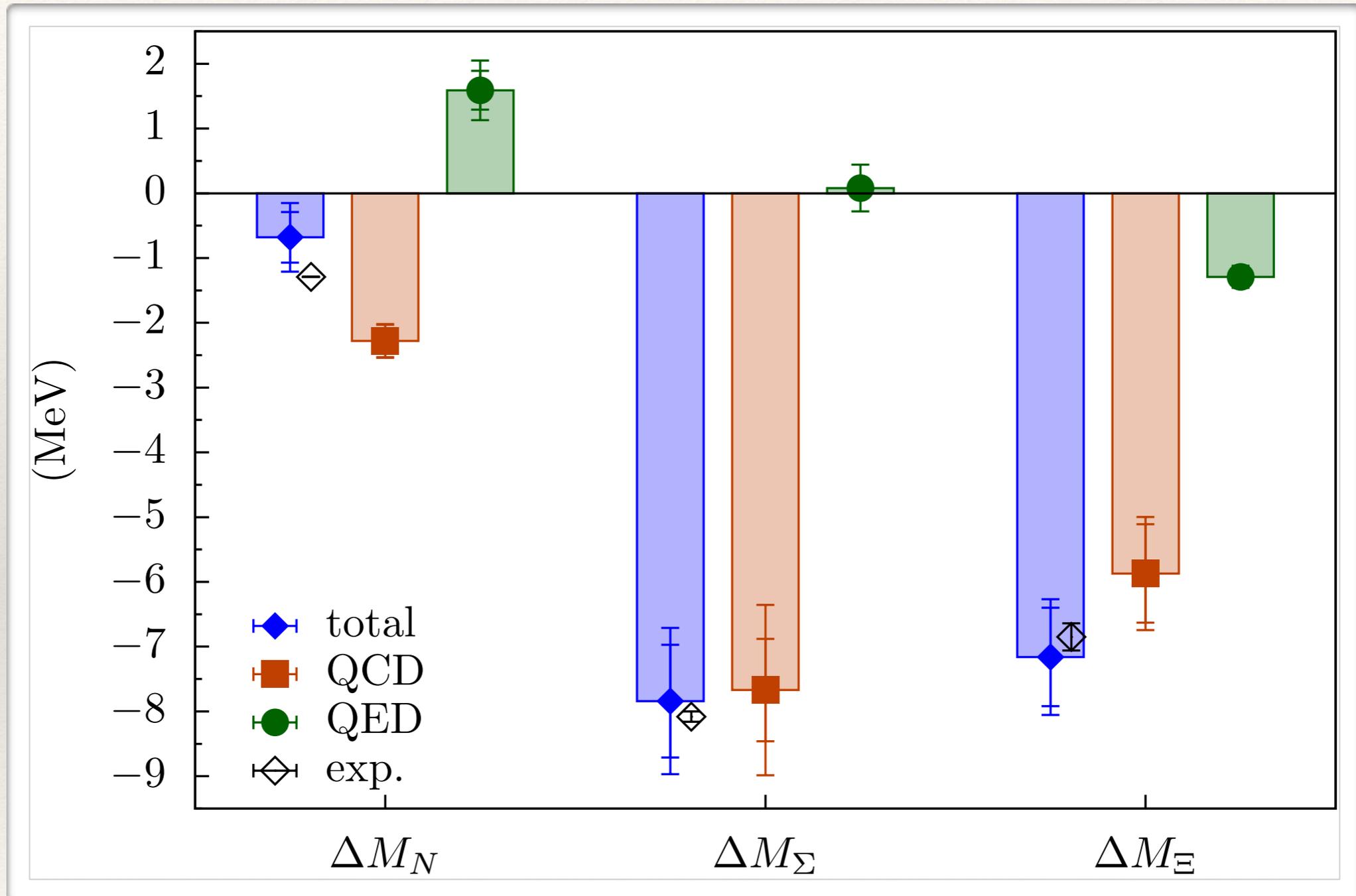
- ❖ How large are the corrections? FLAG parametrisation:

$$\varepsilon = \frac{\Delta_{\text{QED}} M_K^2 - \Delta_{\text{QED}} M_\pi^2}{\Delta M_\pi^2}$$

- ❖ ε is important to determine light quark mass ratios

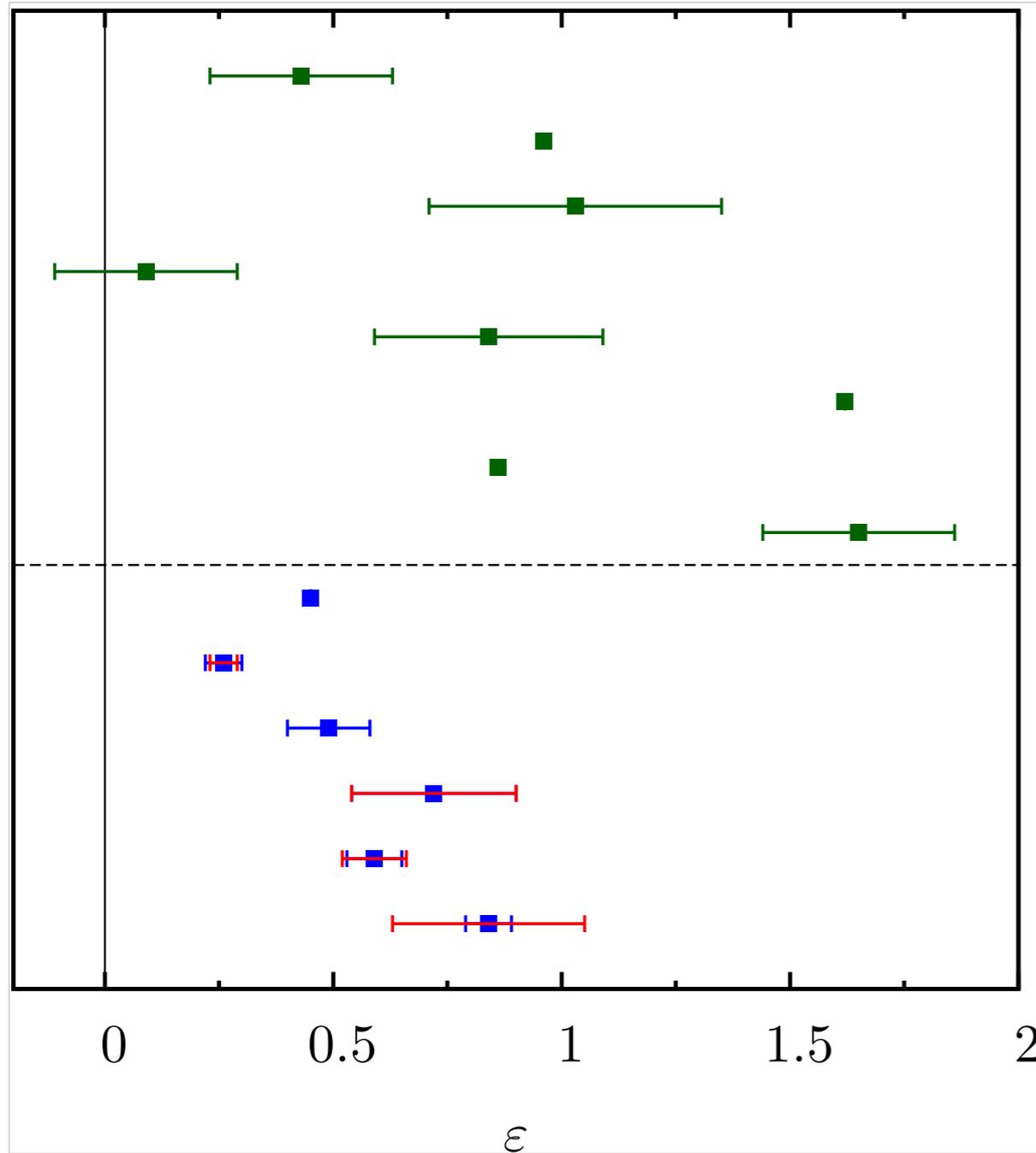
Update on electro-quenched results

EQ results for the baryon spectrum



[BMWc, 2013] (EQ): PRL 111(25) 252001, arXiv: 1306.2287

EQ results for ε



- [Maltman and Kotchan, 1990]
- [Donoghue *et al.*, 1993]
- [Bijnens, 1993]
- [Baur and Urech, 1996]
- [Bijnens and Prades, 1997]
- [Donoghue and Perez, 1997]
- [Gao *et al.*, 1997]
- [Moussallam, 1997]
- [Duncan *et al.*, 1996] (quenched QCD)
- [RBC-UKQCD, 2007]
- [RBC-UKQCD, 2010]
- [RM123, 2013]
- [BMWc, Q, 2014] (preliminary)
- [MILC, 2014] (preliminary)

EQ results for light quark masses

**F. Sanfilippo plenary talk on quark masses:
right after this talk**

Lattice QED

Non-compact lattice QED

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

Non-compact lattice QED

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

- ❖ Pure gauge theory is **free**, it can be solved **exactly**

Non-compact lattice QED

- ❖ Naively discretised **Maxwell action**:

$$S[A_\mu] = \frac{1}{4} \sum_{\mu, \nu} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2$$

- ❖ Pure gauge theory is **free**, it can be solved **exactly**
- ❖ Gauge invariance is preserved

Zero-mode subtraction

Finite volume: **momentum quantisation**

$$\alpha \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2} \cdots \quad \mapsto \quad \frac{\alpha}{V} \sum_k \frac{1}{k^2} \cdots$$

Zero-mode subtraction

Finite volume: **momentum quantisation**

$$\alpha \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2} \cdots \quad \mapsto \quad \frac{\alpha}{V} \sum_k \frac{1}{k^2} \cdots$$



Possibly IR divergent, but
not for physical quantities

Zero-mode subtraction

Finite volume: **momentum quantisation**

$$\alpha \int \frac{d^4 k}{(2\pi)^4} \frac{1}{k^2} \cdots \quad \longrightarrow \quad \frac{\alpha}{V} \sum_k \frac{1}{k^2} \cdots$$



Possibly IR divergent, but
not for physical quantities



Contains a straight $1/0$!

Zero-mode subtraction

- ❖ This problem can be solved by **removing zero modes**

Zero-mode subtraction

- ❖ This problem can be solved by **removing zero modes**
- ❖ **Many possible schemes:**
modification of $A_\mu(k)$ on a set of measure 0

Zero-mode subtraction

- ❖ This problem can be solved by **removing zero modes**
- ❖ **Many possible schemes:**
modification of $A_\mu(k)$ on a set of measure 0
- ❖ Different schemes: **different finite volume behaviours**

Zero-mode subtraction

- ❖ This problem can be solved by **removing zero modes**
- ❖ **Many possible schemes:**
modification of $A_\mu(k)$ on a set of measure 0
- ❖ Different schemes: **different finite volume behaviours**
- ❖ Some more interesting than others

QED_{TL} zero-mode subtraction

❖ QED_{TL}: $A_\mu(0) = 0$

Mostly used in all simulations so far

QED_{TL} zero-mode subtraction

❖ QED_{TL}: $A_\mu(0) = 0$

Mostly used in all simulations so far

❖ With QED_{TL}, the $T \rightarrow \infty$, $L = \text{cst.}$ limit **can diverge**:

$$\frac{\alpha}{V} \sum_{k \neq 0} \frac{1}{k^2} \cdots \quad \longmapsto \quad \frac{\alpha}{L^3} \int \frac{dk_0}{2\pi} \sum_{\mathbf{k}} \frac{1}{k^2} \cdots$$

QED_{TL} zero-mode subtraction

❖ QED_{TL}: $A_\mu(0) = 0$

Mostly used in all simulations so far

❖ With QED_{TL}, the $T \rightarrow \infty$, $L = \text{cst.}$ limit **can diverge**:

$$\frac{\alpha}{V} \sum_{k \neq 0} \frac{1}{k^2} \cdots \quad \longmapsto \quad \frac{\alpha}{L^3} \int \frac{dk_0}{2\pi} \sum_{\mathbf{k}} \frac{1}{k^2} \cdots$$

❖ QED_{TL} **does not have reflection positivity**

QED_{TL} finite-volume effects

❖ Example — 1-loop QED_{TL} [BMWc, 2014]:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \left[1 - \frac{\pi T}{2\kappa L} \right] \right) - \frac{3\pi}{(mL)^3} \left[1 - \frac{\coth(mT)}{2} \right] - \frac{3\pi L}{2(mL)^4 T} \right] \right\}$$

up to exponential corrections, with $\kappa = 2.83729 \dots$

QED_{TL} finite-volume effects

- ❖ Example — 1-loop QED_{TL} [BMWc, 2014]:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \left[1 - \frac{\pi T}{2\kappa L} \right] \right) - \frac{3\pi}{(mL)^3} \left[1 - \frac{\coth(mT)}{2} \right] - \frac{3\pi L}{2(mL)^4 T} \right] \right\}$$

up to exponential corrections, with $\kappa = 2.83729 \dots$

- ❖ **Divergent finite volume effects** with $T \rightarrow \infty$, $L = \text{cst.}$

QED_{TL} finite-volume effects

- ❖ Example — 1-loop QED_{TL} [BMWc, 2014]:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \left[1 - \frac{\pi T}{2\kappa L} \right] \right) - \frac{3\pi}{(mL)^3} \left[1 - \frac{\coth(mT)}{2} \right] - \frac{3\pi L}{2(mL)^4 T} \right] \right\}$$

up to exponential corrections, with $\kappa = 2.83729 \dots$

- ❖ **Divergent finite volume effects** with $T \rightarrow \infty$, $L = \text{cst.}$
- ❖ Same behaviour independently discovered by MILC

QED_L zero-mode subtraction

- ❖ QED_L: $A_\mu(k_0, \mathbf{0}) = 0$
inspired from [Hayakawa & Uno, 2008]

QED_L zero-mode subtraction

- ❖ QED_L: $A_\mu(k_0, \mathbf{0}) = 0$
inspired from [Hayakawa & Uno, 2008]
- ❖ QED_L maintains reflection positivity [BMWc, 2014]:

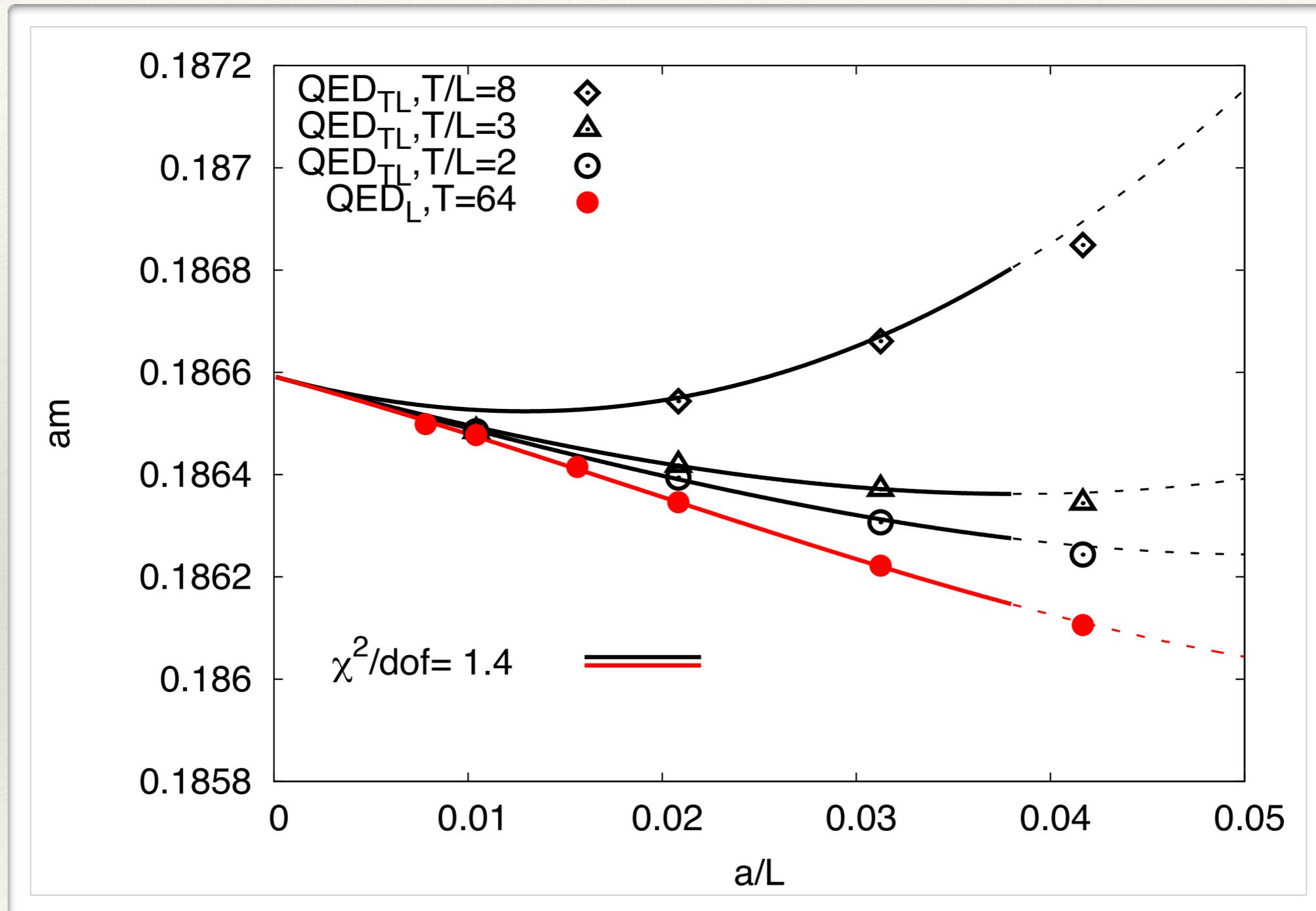
QED_L zero-mode subtraction

- ❖ QED_L: $A_\mu(k_0, \mathbf{0}) = 0$
inspired from [Hayakawa & Uno, 2008]
- ❖ QED_L maintains reflection positivity [BMWc, 2014]:
- ❖ QED_L finite volume effects:

$$m(T, L) \underset{T, L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) - \frac{3\pi}{(mL)^3} \right] \right\}$$

inverse powers of L, independent of T

Finite-volume effects



Pure QED simulations (quenched) from [BMWc, 2014]

Finite-volume effects

- ❖ What about **composite particles** (QCD + QED)?

Finite-volume effects

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT

Finite-volume effects

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons

Finite-volume effects

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons
- ❖ [Davoudi & Savage, 2014]: NREFTs
mesons, baryons, nuclei and HVP

$$m(L) \underset{L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) + \mathcal{O} \left(\frac{1}{L^3} \right) \right] \right\}$$

Finite-volume effects

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons
- ❖ [Davoudi & Savage, 2014]: NREFTs
mesons, baryons, nuclei and HVP

$$m(L) \underset{L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) + \mathcal{O} \left(\frac{1}{L^3} \right) \right] \right\}$$

- ❖ [BMWc, 2014]: Ward identities: **NLO is universal**

Finite-volume effects

- ❖ What about **composite particles** (QCD + QED)?
- ❖ [Hayakawa & Uno, 2008]: SU(3) PQChPT
- ❖ [RBC-UKQCD, 2010]: SU(2) PQChPT + heavy kaons
- ❖ [Davoudi & Savage, 2014]: NREFTs
mesons, baryons, nuclei and HVP

$$m(L) \underset{L \rightarrow +\infty}{\sim} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) + \mathcal{O} \left(\frac{1}{L^3} \right) \right] \right\}$$

- ❖ [BMWc, 2014]: Ward identities: **NLO is universal**
- ❖ parallel talk by **C. Lehner: tomorrow 15:35**

Full QCD+QED simulations

Full QCD + QED projects

	RBC-UKQCD	PACS-CS	QCDSF-UKQCD	BMW _c
arXiv	1006.1311	1205.2961	1311.4554 and Lat. 2014	1406.4088
fermions	DWF	clover	clover	clover
N_f	2+1	1+1+1	1+1+1	1+1+1+1
method	reweighting	reweighting	RHMC	RHMC
$\min(M_\pi)$ (MeV)	420	135	250	195
a (fm)	0.11	0.09	0.08	0.06 — 0.10
$\#a$	1	1	1	4
L (fm)	1.8	2.9	1.9 — 2.6	2.1 — 8.3
$\#L$	1	1	2	11

Starting simulation program by MILC: **R. Zhou talk Monday 14:15**

[BMWc, 2014]: QED simulations

- ❖ No mass gap: **large autocorrelations !**

[BMWc, 2014]: QED simulations

- ❖ No mass gap: **large autocorrelations !**
- ❖ One can determine exactly an MD Hamiltonian that **removes all memory** in the QED Markov chain:

$$H = \frac{1}{2TL^3} \sum_{\mu, k} \left\{ |\hat{k}|^2 |A_{\mu, k}|^2 + \frac{\pi}{4|\hat{k}|^2} |\Pi_{\mu, k}|^2 \right\}$$

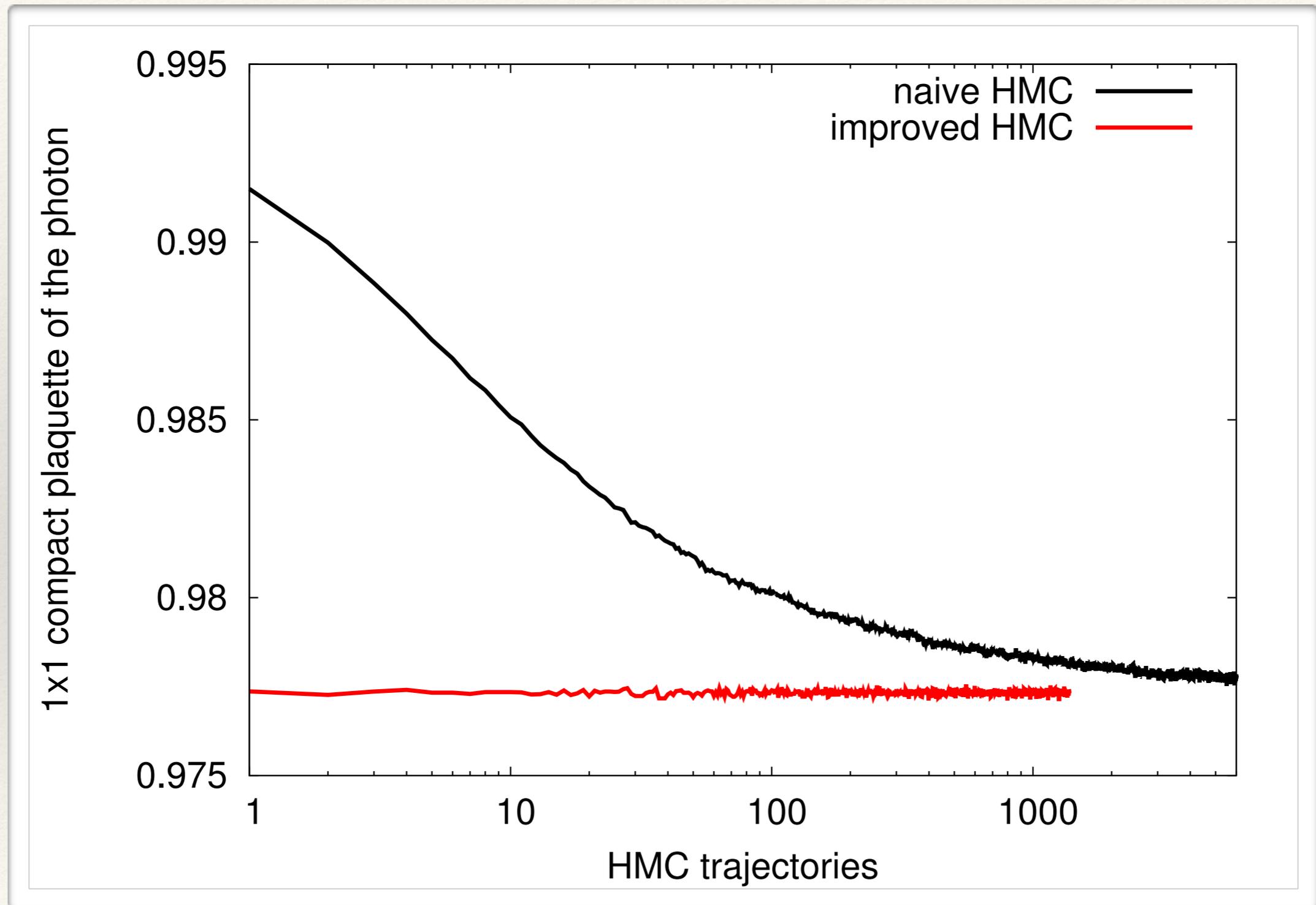
[BMWc, 2014]: QED simulations

- ❖ No mass gap: **large autocorrelations !**
- ❖ One can determine exactly an MD Hamiltonian that **removes all memory** in the QED Markov chain:

$$H = \frac{1}{2TL^3} \sum_{\mu, k} \left\{ |\hat{k}|^2 |A_{\mu, k}|^2 + \frac{\pi}{4|\hat{k}|^2} |\Pi_{\mu, k}|^2 \right\}$$

- ❖ Clover term greatly reduces discretisation errors

[BMWc, 2014]: QED simulations



Isospin splittings
in the hadron spectrum

[QCDSF, 2014]: progress summary

[QCDSF, 2014]: progress summary

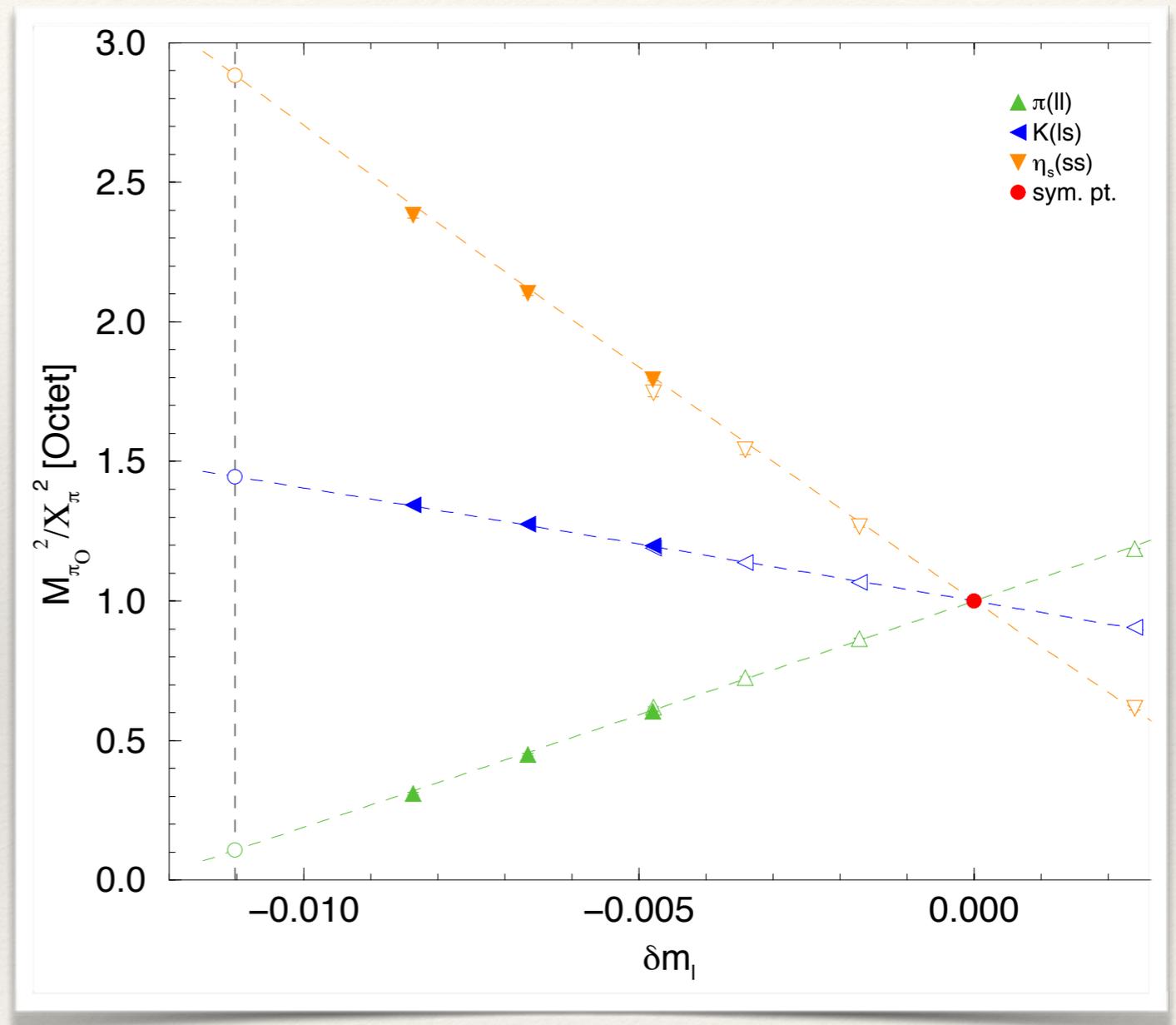
- ❖ $N_f = 1+1+1$
full QCD+QED
simulations in progress

[QCDSF, 2014]: progress summary

- ❖ $N_f = 1+1+1$
full QCD+QED
simulations in progress
- ❖ computational strategy
is a continuation of
[arXiv:1102.5300]

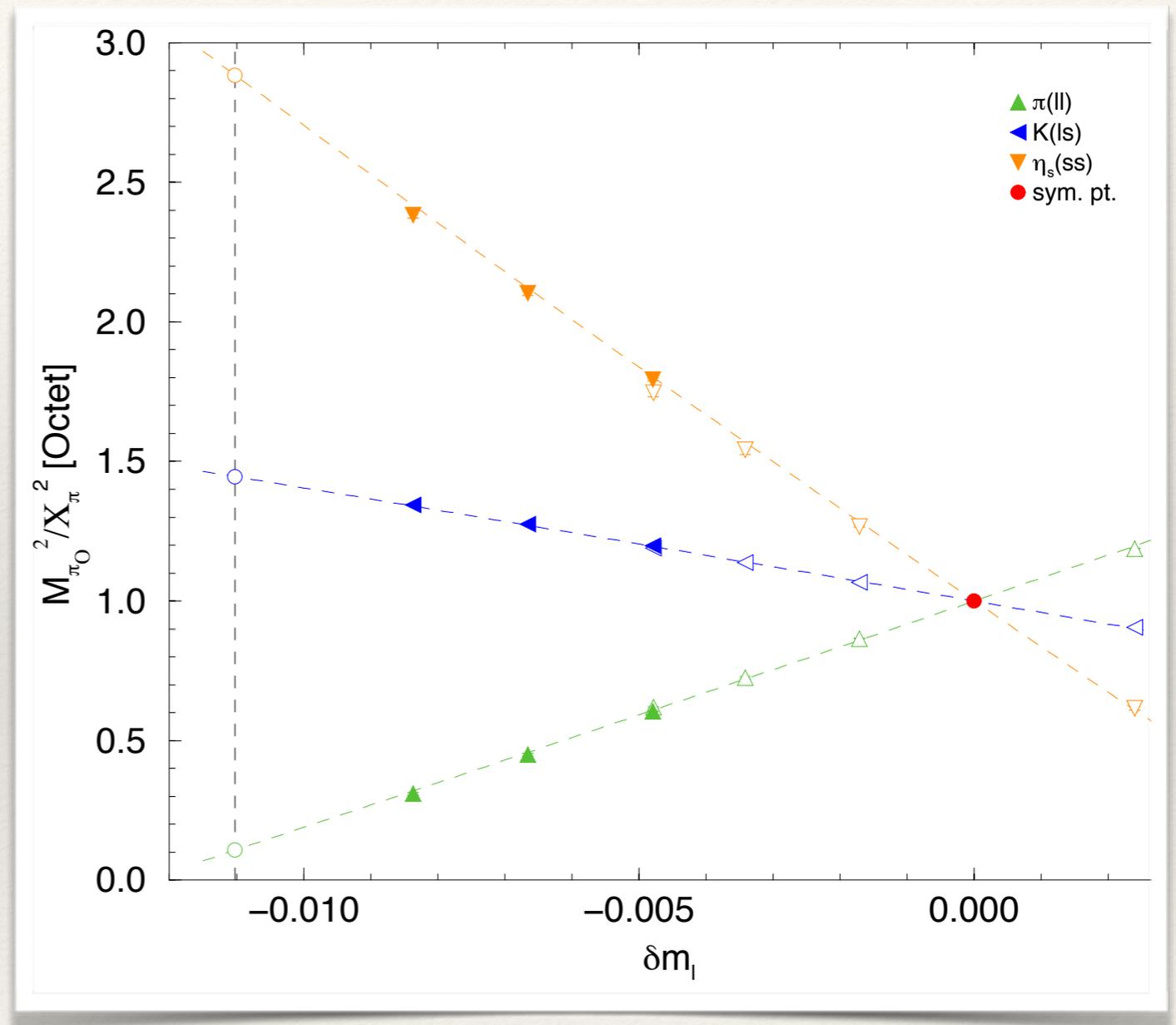
[QCDSF, 2014]: progress summary

- ❖ $N_f = 1+1+1$
full QCD+QED
simulations in progress
- ❖ computational strategy
is a continuation of
[arXiv:1102.5300]
- ❖ start from the SU(3)
symmetric point and
move keeping
 $m_u + m_d + m_s$
constant



[QCDSF, 2014]: progress summary

- ❖ $N_f = 1+1+1$
full QCD+QED
simulations in progress
- ❖ computational strategy
is a continuation of
[arXiv:1102.5300]
- ❖ start from the SU(3)
symmetric point and
move keeping
 $m_u + m_d + m_s$
constant



G. Schierholz parallel talk: tomorrow 14:15

[BMWc, 2014]: mass splitting calculation

[BMWc, 2014]: mass splitting calculation

- ❖ **many smeared sources per configurations ($O(100)$)**

[BMWc, 2014]: mass splitting calculation

- ❖ **many smeared sources** per configurations ($O(100)$)
- ❖ electric charge renormalisation using **Wilson flow**

[BMWc, 2014]: mass splitting calculation

- ❖ **many smeared sources** per configurations ($O(100)$)
- ❖ electric charge renormalisation using **Wilson flow**
- ❖ small extrapolation to the physical point
(similar to [BMWc, 2013])

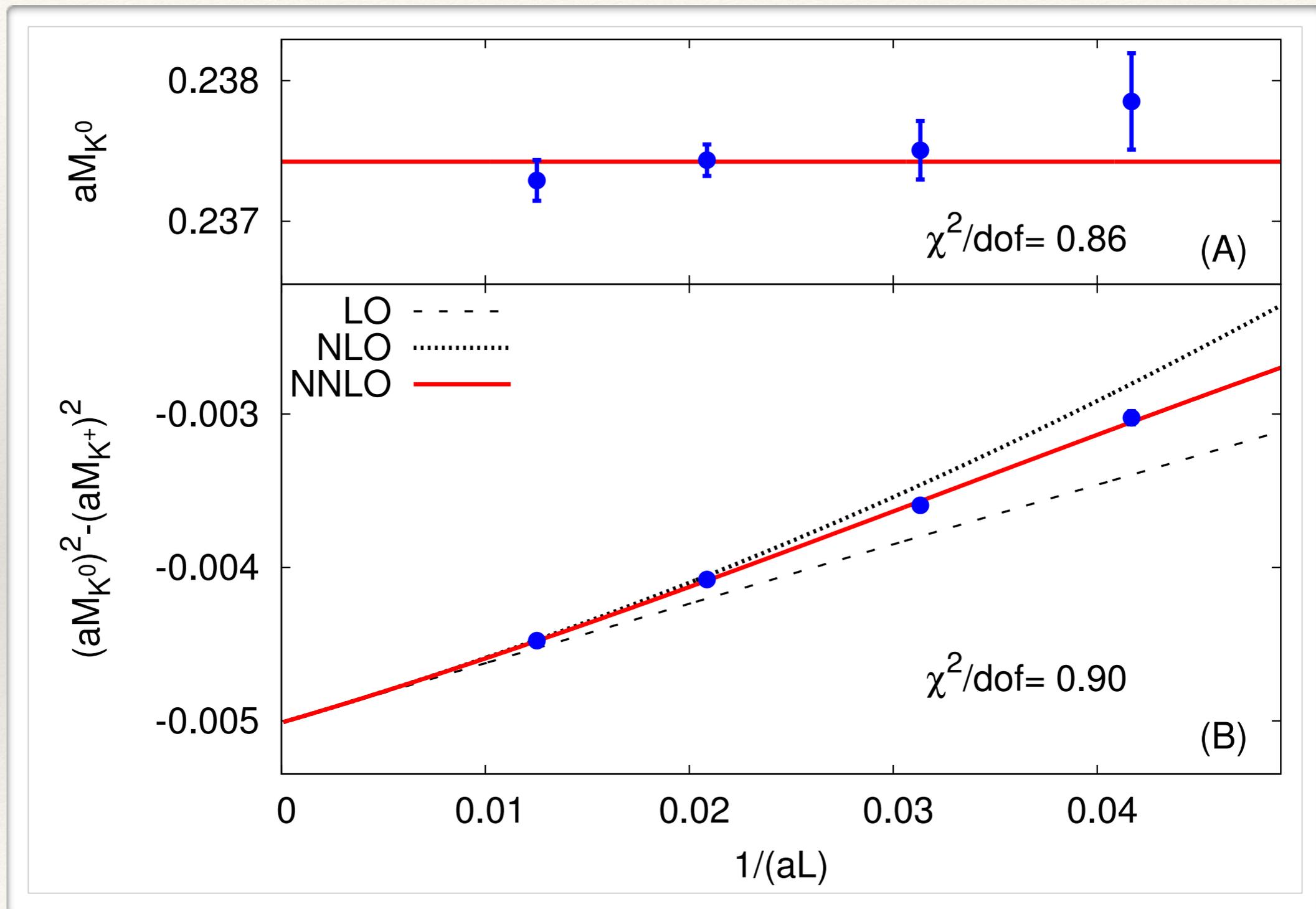
[BMWc, 2014]: mass splitting calculation

- ❖ **many smeared sources** per configurations ($O(100)$)
- ❖ electric charge renormalisation using **Wilson flow**
- ❖ small extrapolation to the physical point
(similar to [BMWc, 2013])
- ❖ Systematic error based on BMW's histogram method.
Weights are based on the goodness of the fits, flat and Akaike's information criterion (**overfitting is penalised**)

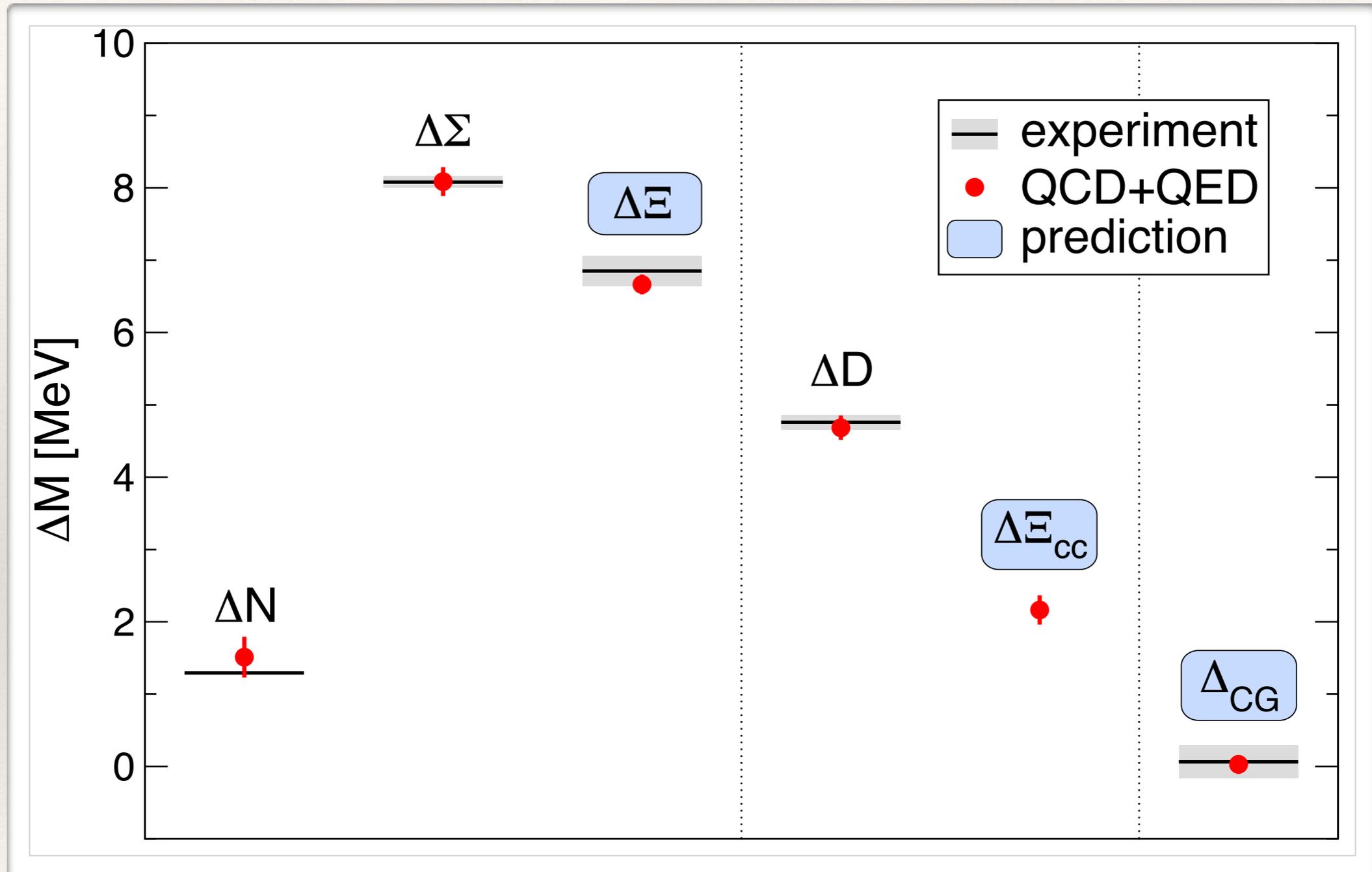
[BMWc, 2014]: mass splitting calculation

- ❖ **many smeared sources** per configurations ($O(100)$)
- ❖ electric charge renormalisation using **Wilson flow**
- ❖ small extrapolation to the physical point
(similar to [BMWc, 2013])
- ❖ Systematic error based on BMW's histogram method.
Weights are based on the goodness of the fits, flat and Akaike's information criterion (**overfitting is penalised**)
- ❖ **$O(500)$ analyses per mass splitting**

[BMWc, 2014]: finite-volume study

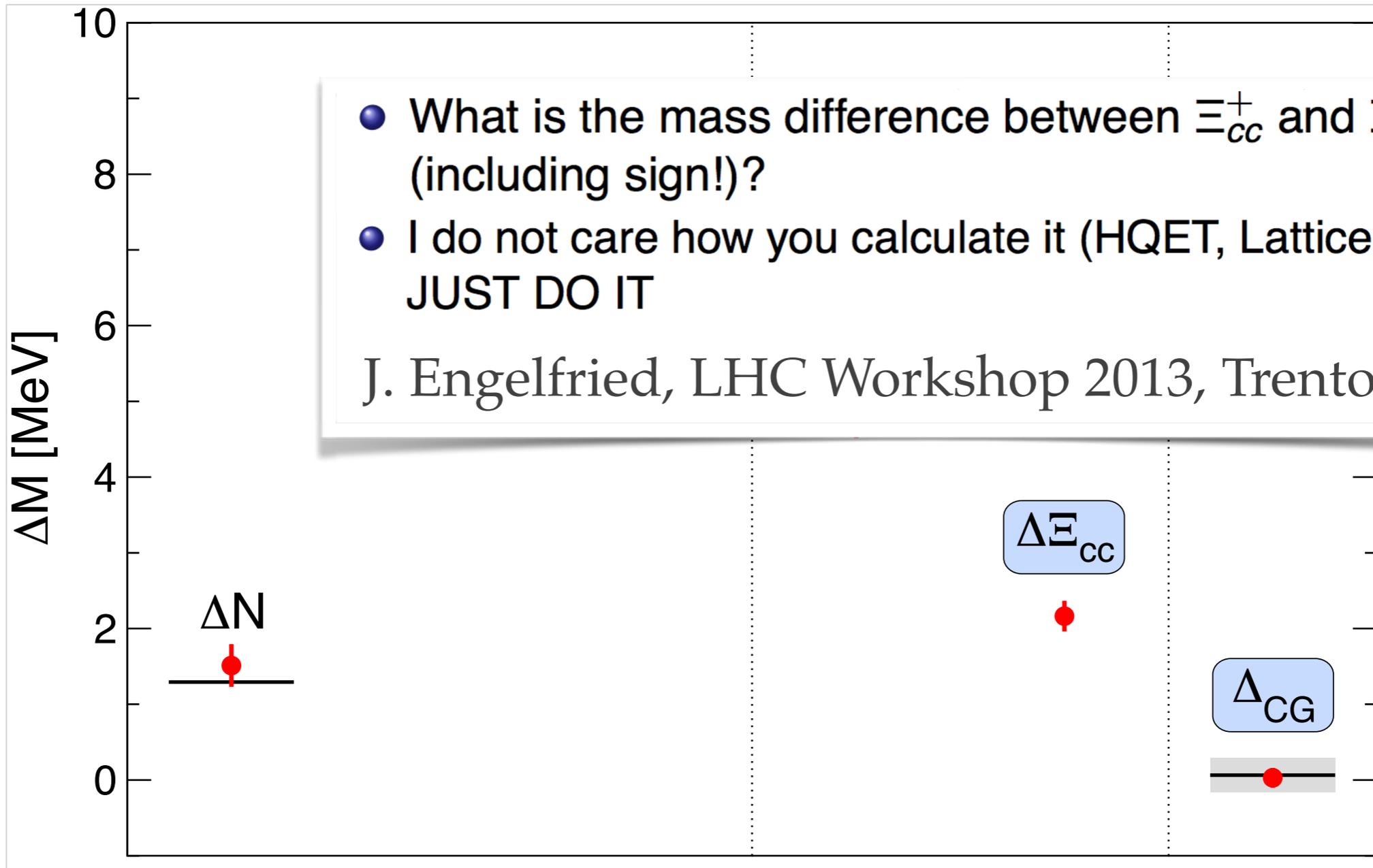


[BMWc, 2014]: result summary



$$\Delta_{CG} = \Delta M_N - \Delta M_\Sigma + \Delta M_\Xi \text{ (Coleman-Glashow relation)}$$

[BMWc, 2014]: result summary

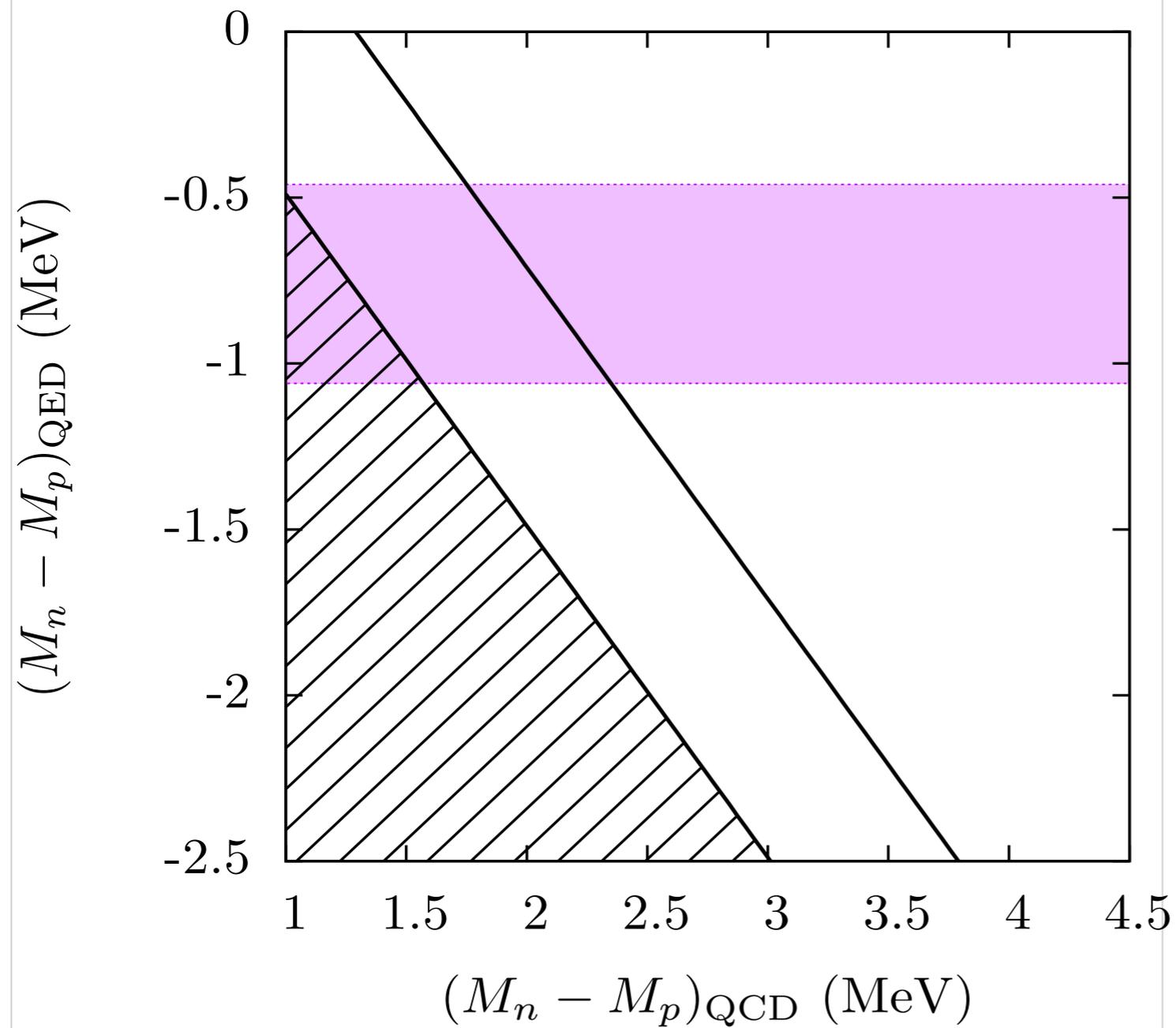


- What is the mass difference between Ξ_{cc}^+ and Ξ_{cc}^{++} (including sign!)?
- I do not care how you calculate it (HQET, Lattice, ...), JUST DO IT

J. Engelfried, LHC Workshop 2013, Trento

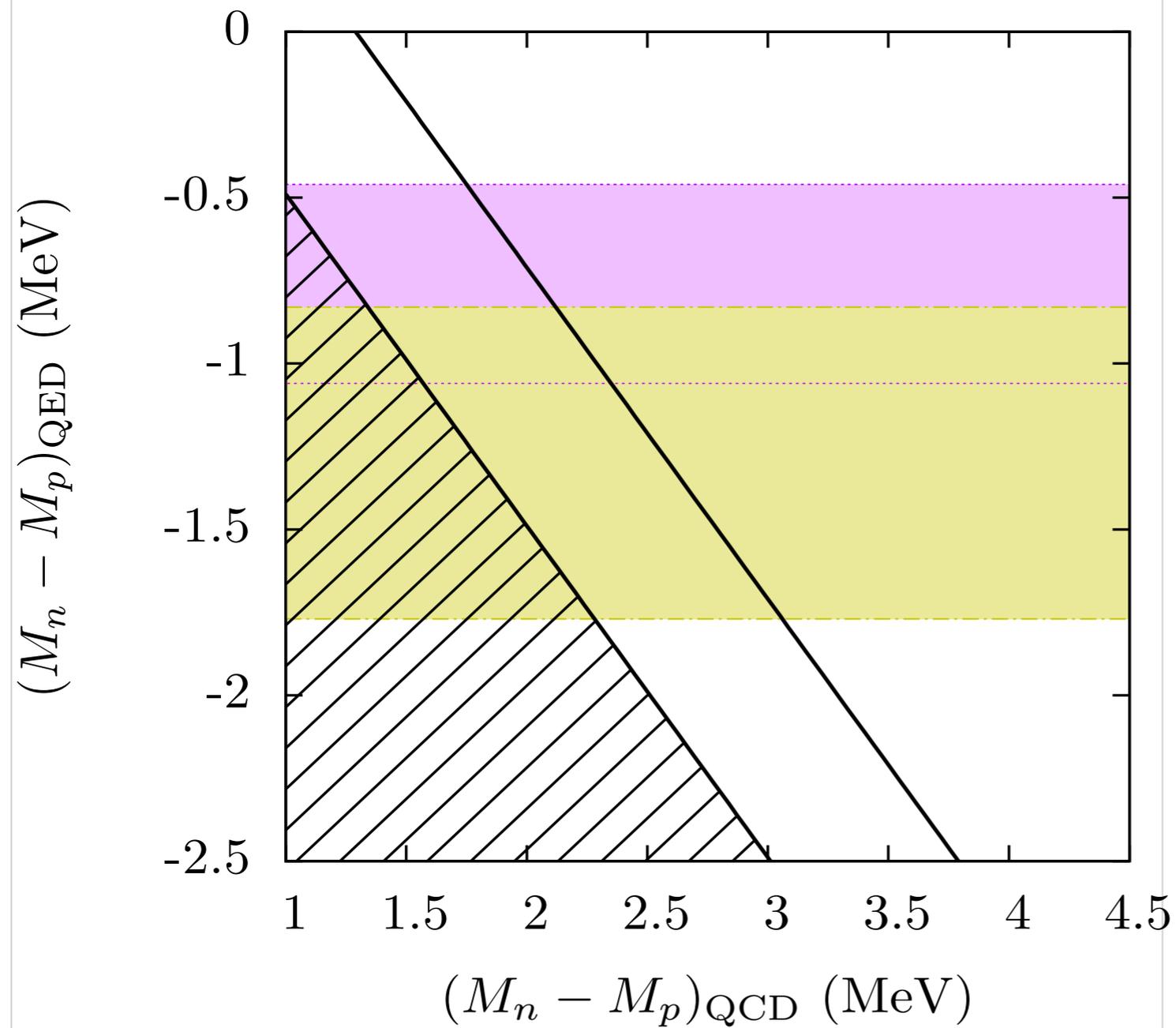
$$\Delta_{CG} = \Delta M_N - \Delta M_\Sigma + \Delta M_\Xi \text{ (Coleman-Glashow relation)}$$

Results for the nucleon mass splitting



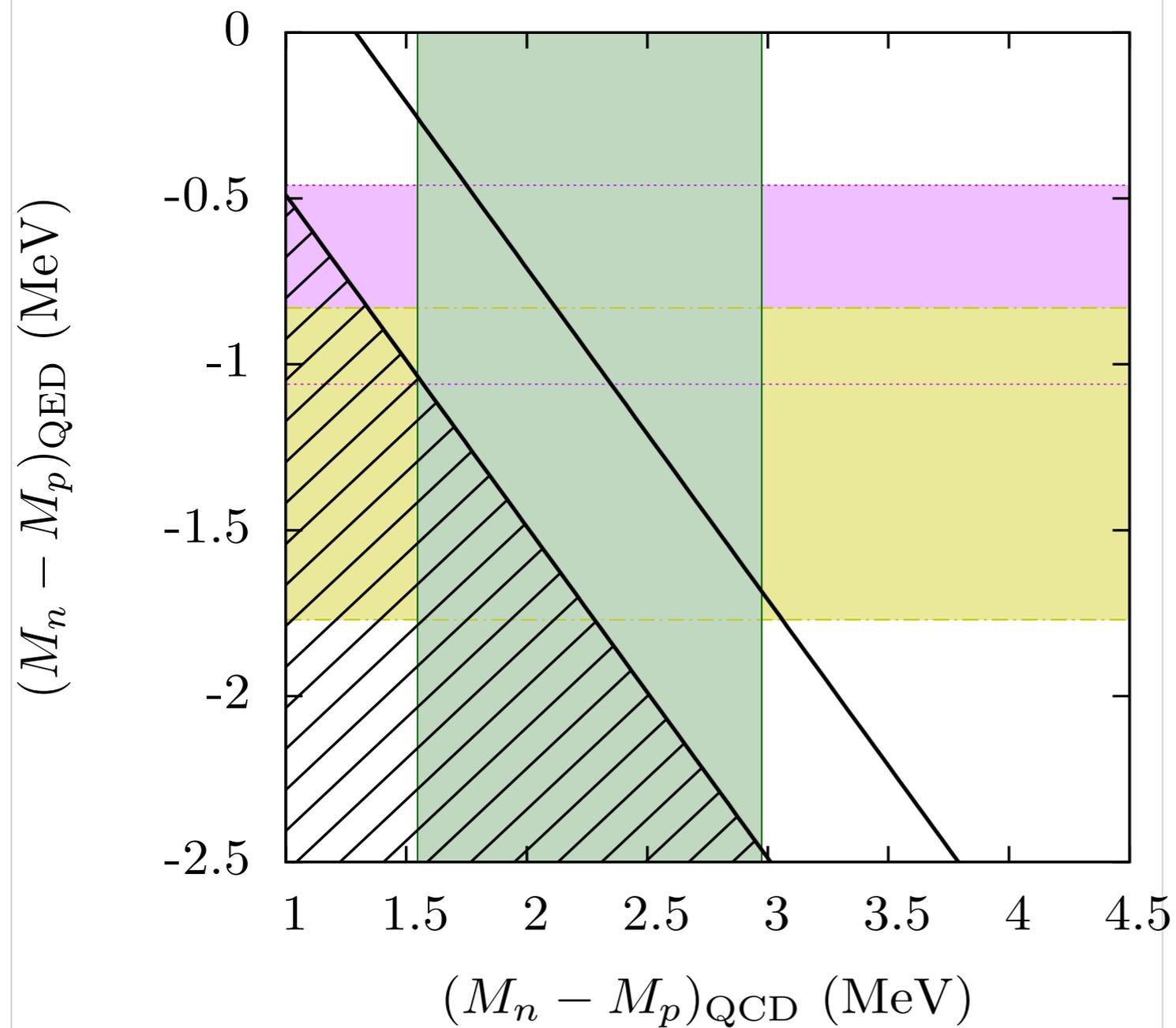
- [Gasser & Leutwyler, 1982]
- no *beta*-decay
- experiment

Results for the nucleon mass splitting



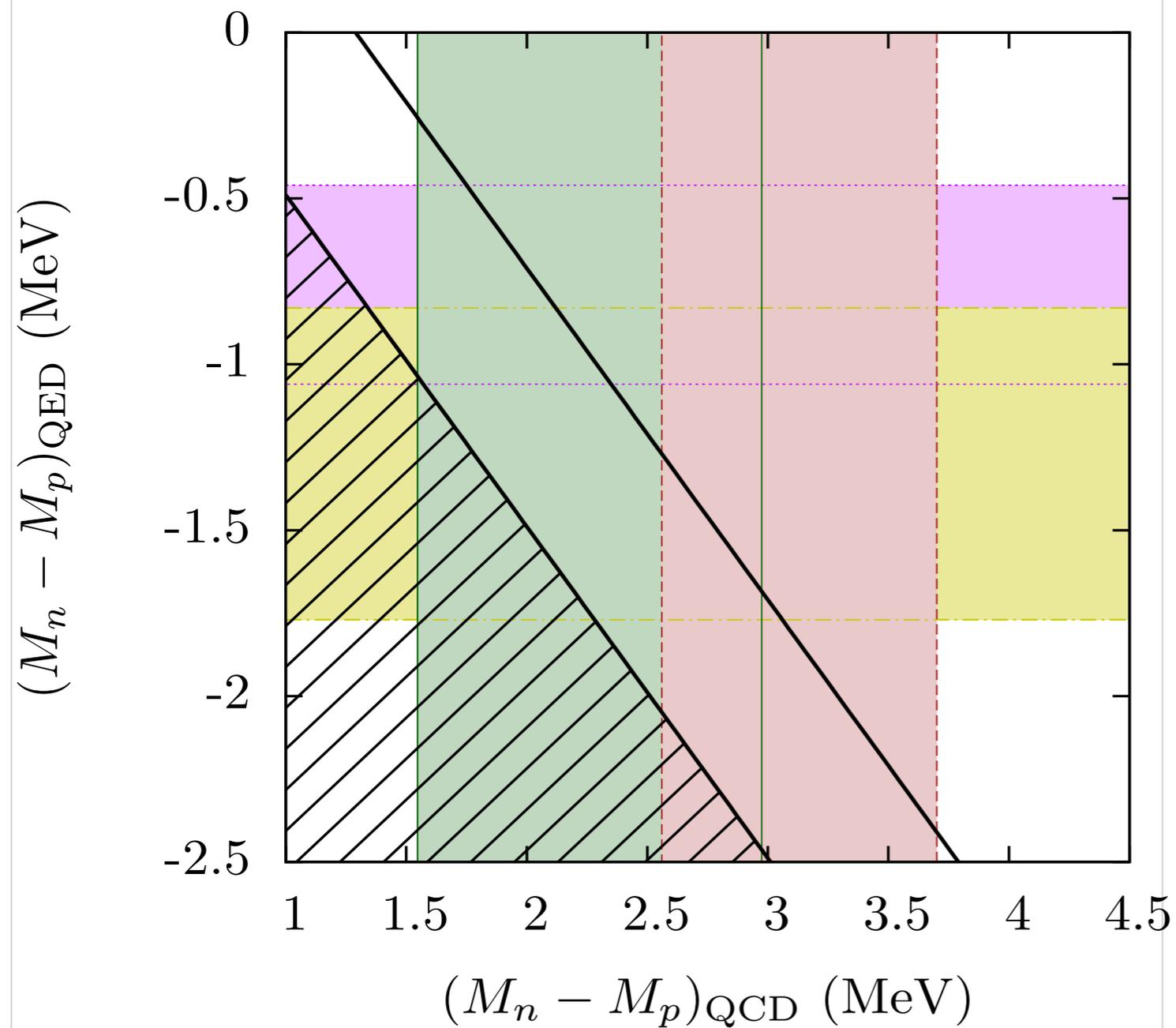
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- no *beta*-decay
- experiment

Results for the nucleon mass splitting



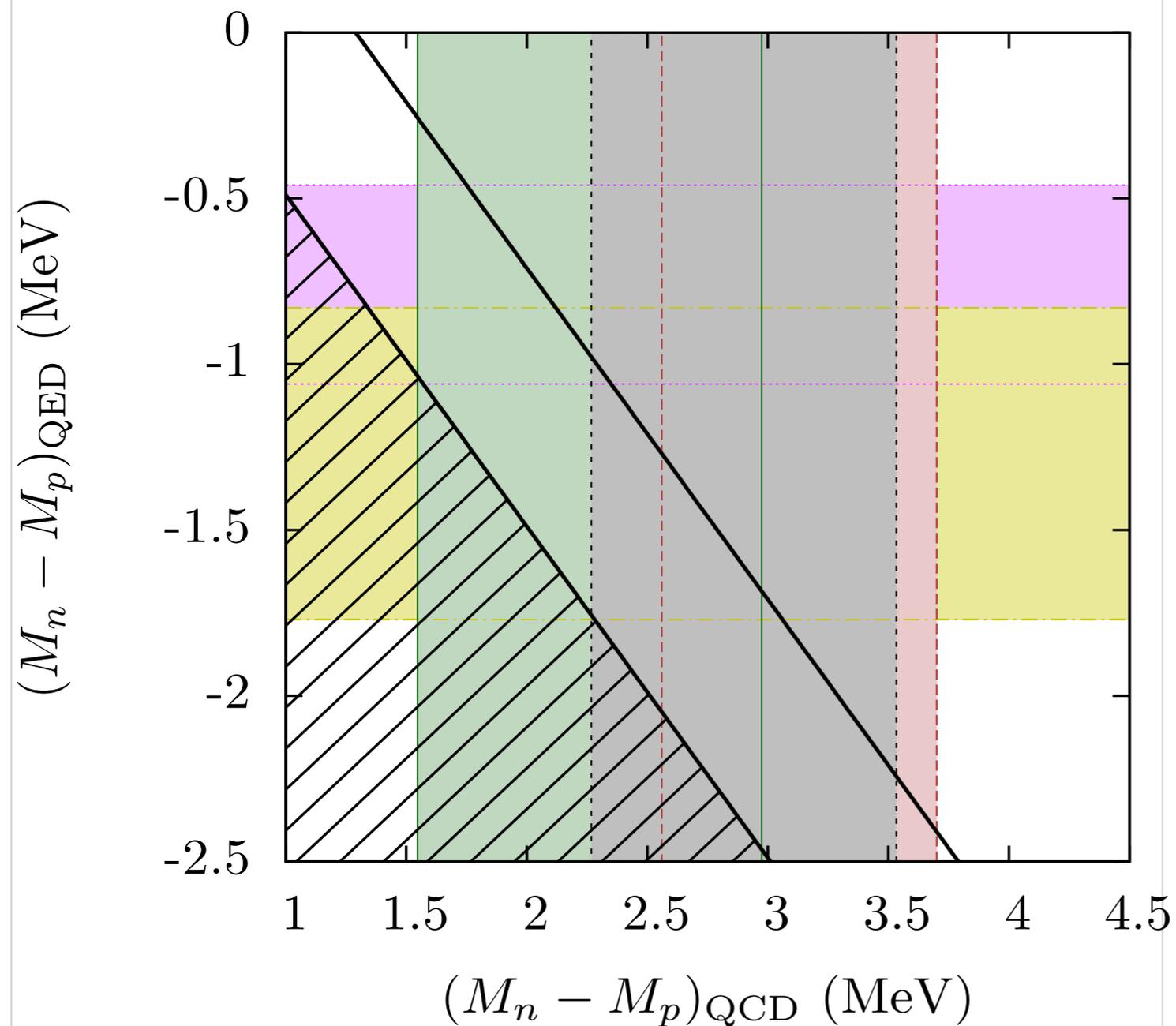
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- no *beta*-decay
- experiment

Results for the nucleon mass splitting



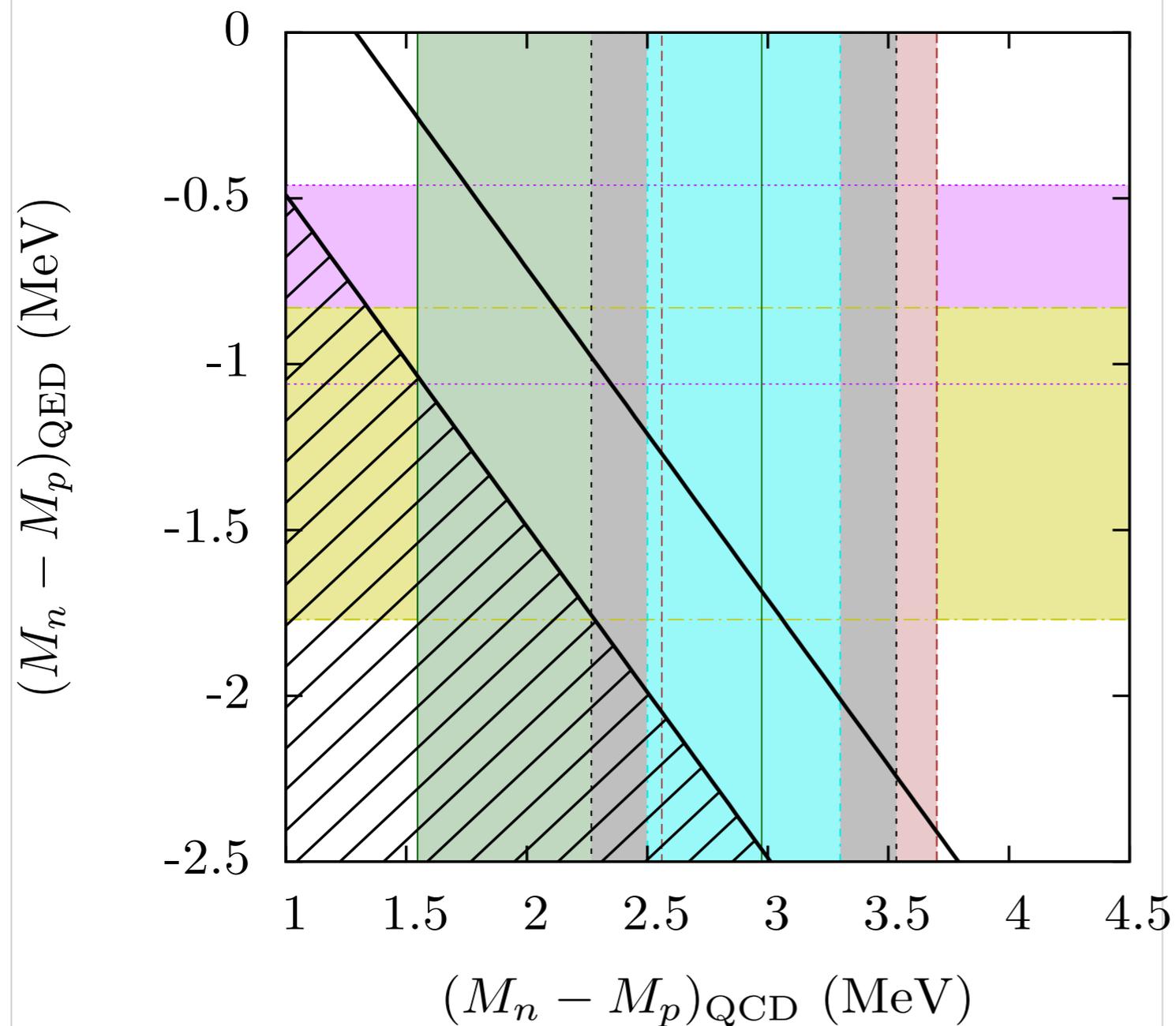
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- no *beta*-decay
- experiment

Results for the nucleon mass splitting



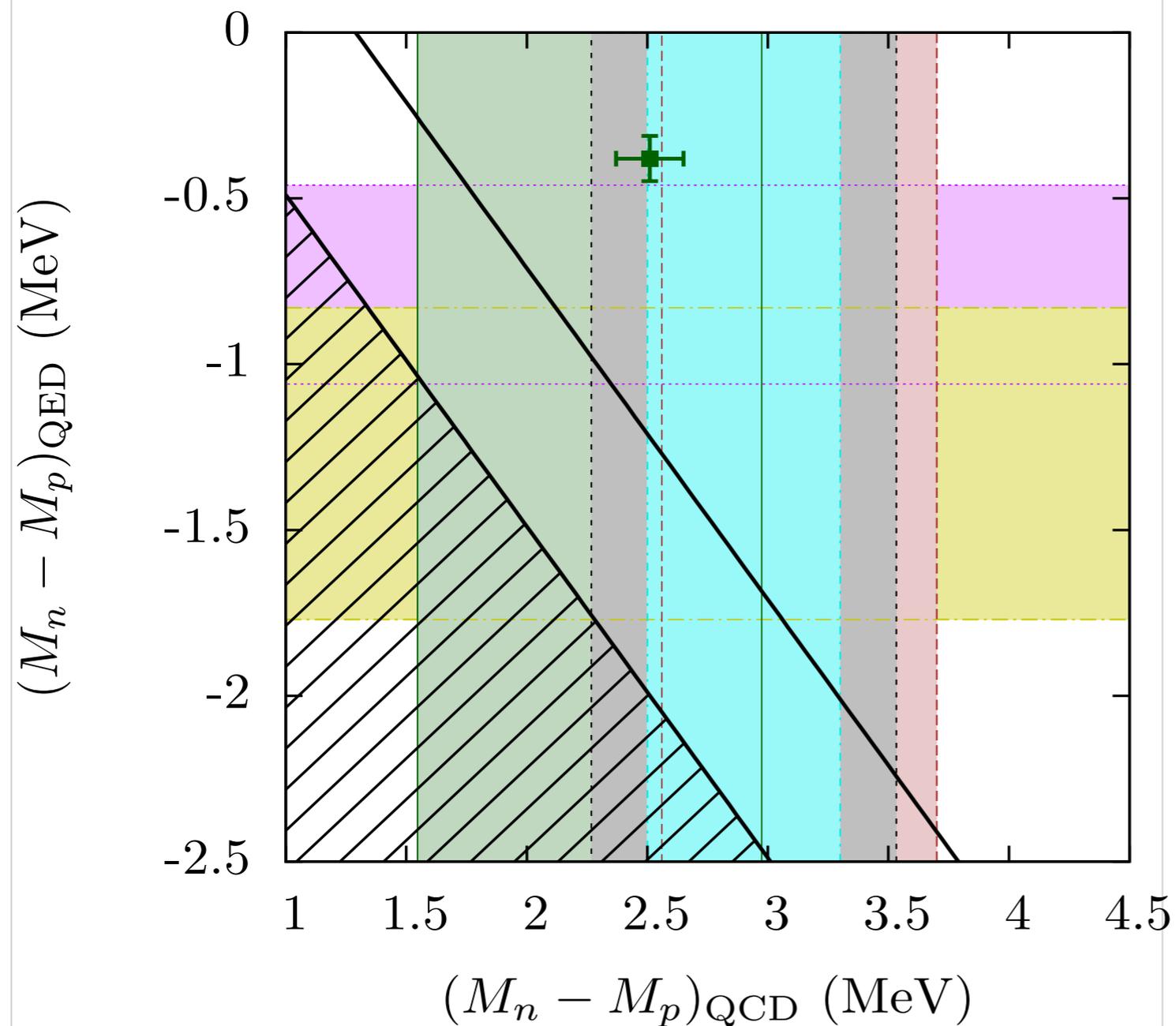
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- no *beta*-decay
- experiment

Results for the nucleon mass splitting



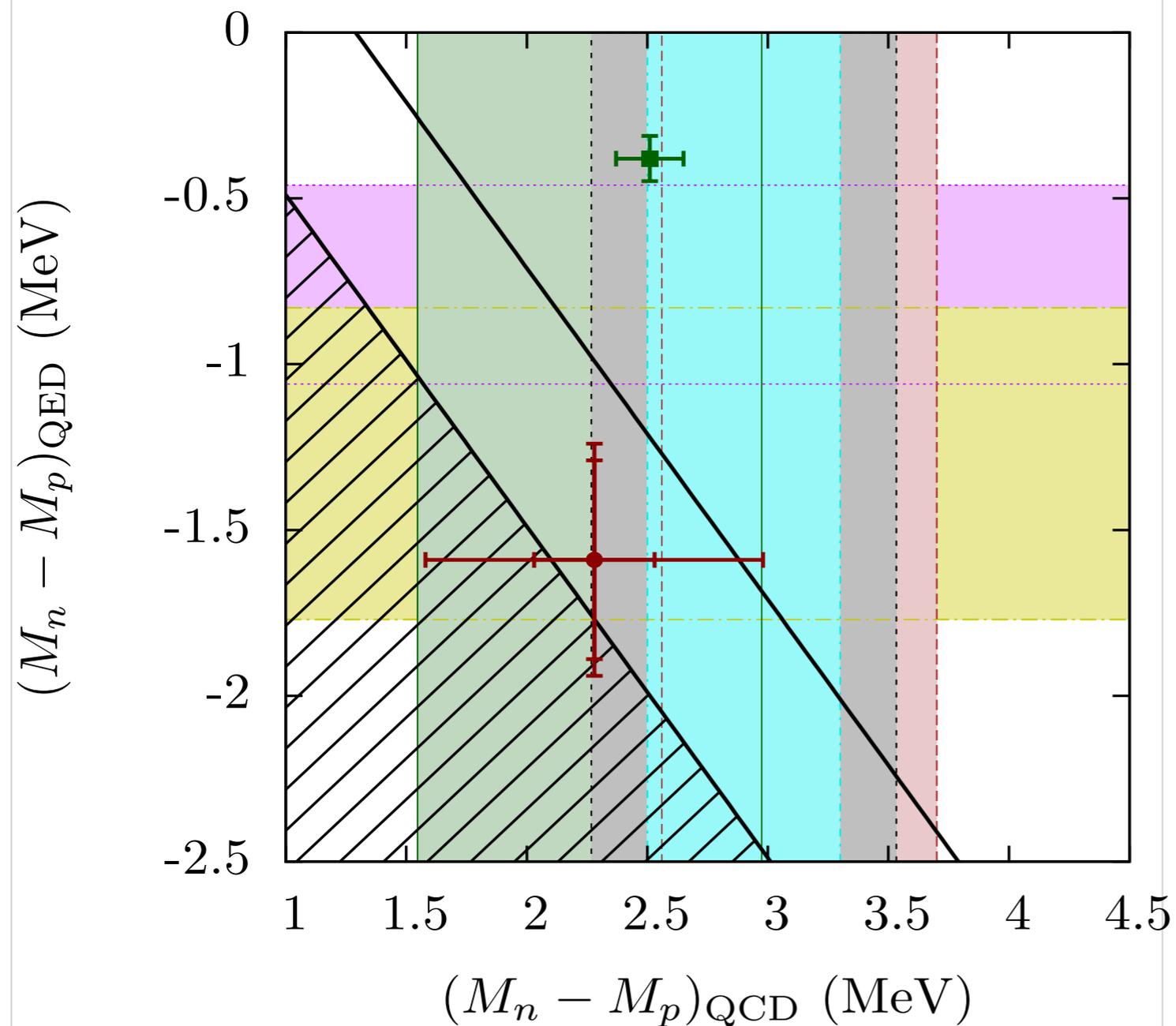
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- [Shanahan *et al.*, 2012]
- no *beta*-decay
- experiment

Results for the nucleon mass splitting



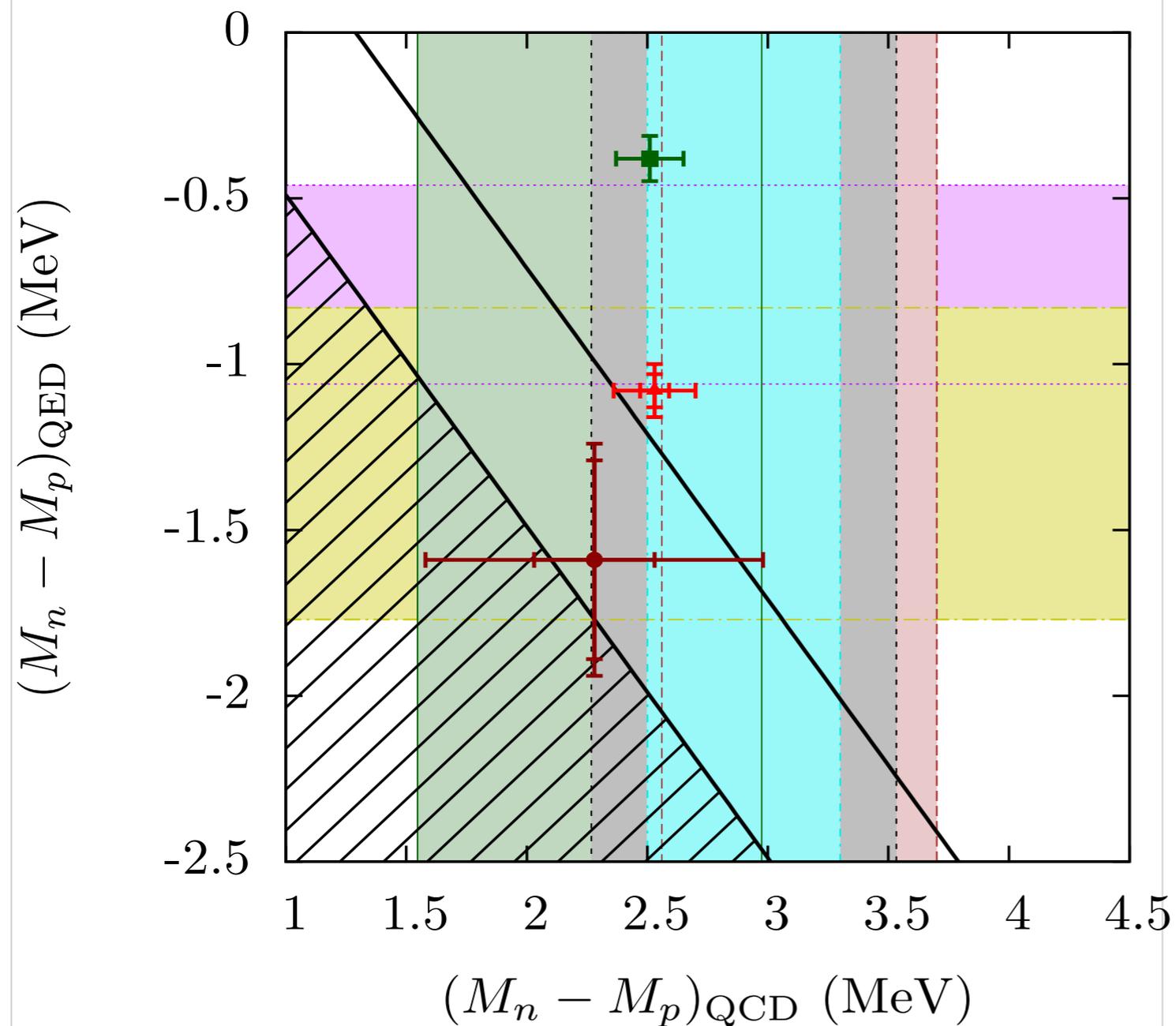
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- [Shanahan *et al.*, 2012]
- no *beta*-decay
- experiment
- [RBC-UKQCD, 2010]

Results for the nucleon mass splitting



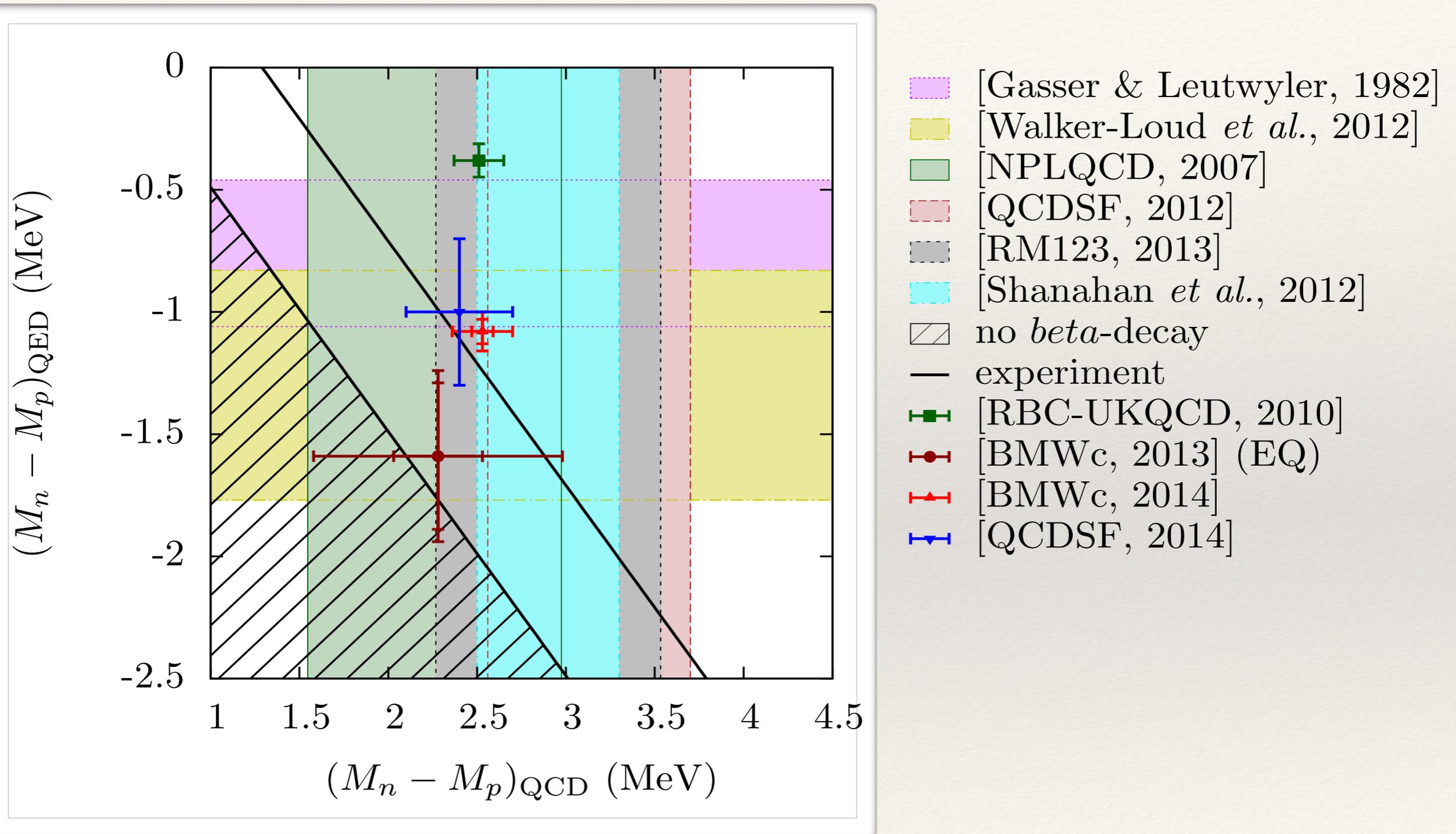
- [Gasser & Leutwyler, 1982]
- [Walker-Loud *et al.*, 2012]
- [NPLQCD, 2007]
- [QCDSF, 2012]
- [RM123, 2013]
- [Shanahan *et al.*, 2012]
- no *beta*-decay
- experiment
- [RBC-UKQCD, 2010]
- [BMWc, 2013] (EQ)

Results for the nucleon mass splitting



- ▨ [Gasser & Leutwyler, 1982]
- ▨ [Walker-Loud *et al.*, 2012]
- ▨ [NPLQCD, 2007]
- ▨ [QCDSF, 2012]
- ▨ [RM123, 2013]
- ▨ [Shanahan *et al.*, 2012]
- ▨ no *beta*-decay
- experiment
- [RBC-UKQCD, 2010]
- [BMWc, 2013] (EQ)
- ▲ [BMWc, 2014]

Results for the nucleon mass splitting



Summary & outlook

Summary

Summary

- ❖ We now have a good understanding of QCD+QED on a finite lattice

Summary

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**

Summary

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**
- ❖ [BMWc, 2014]: **full simulations of the low-energy SM** with a potential precision of $O[(N_c m_b^2)^{-1}, \alpha^2] \sim 10^{-4}$

Summary

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**
- ❖ [BMWc, 2014]: **full simulations of the low-energy SM** with a potential precision of $O[(N_c m_b^2)^{-1}, \alpha^2] \sim 10^{-4}$
- ❖ The isospin splittings in the hadron spectrum are determined with a **high accuracy and full control of uncertainties**

Summary

- ❖ We now have a good understanding of QCD+QED on a finite lattice
- ❖ Finite-size effects on masses are now **well controlled**
- ❖ [BMWc, 2014]: **full simulations of the low-energy SM** with a potential precision of $O[(N_c m_b^2)^{-1}, \alpha^2] \sim 10^{-4}$
- ❖ The isospin splittings in the hadron spectrum are determined with a **high accuracy and full control of uncertainties**
- ❖ The nucleon mass splitting is determined as a $> 5\sigma$ effect

Outlook

Outlook

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections

Outlook

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections
- ❖ QCD+QED decay constants are gauge variant and IR divergent. How to deal with that?
C.T. Sachrajda plenary talk: 28/06 — 10:30

Outlook

- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections
- ❖ QCD+QED decay constants are gauge variant and IR divergent. How to deal with that?
C.T. Sachrajda plenary talk: 28/06 — 10:30
- ❖ Compute corrections to matrix elements
($K_{\ell 3}$, $K \rightarrow \pi\pi, \dots$)

Outlook

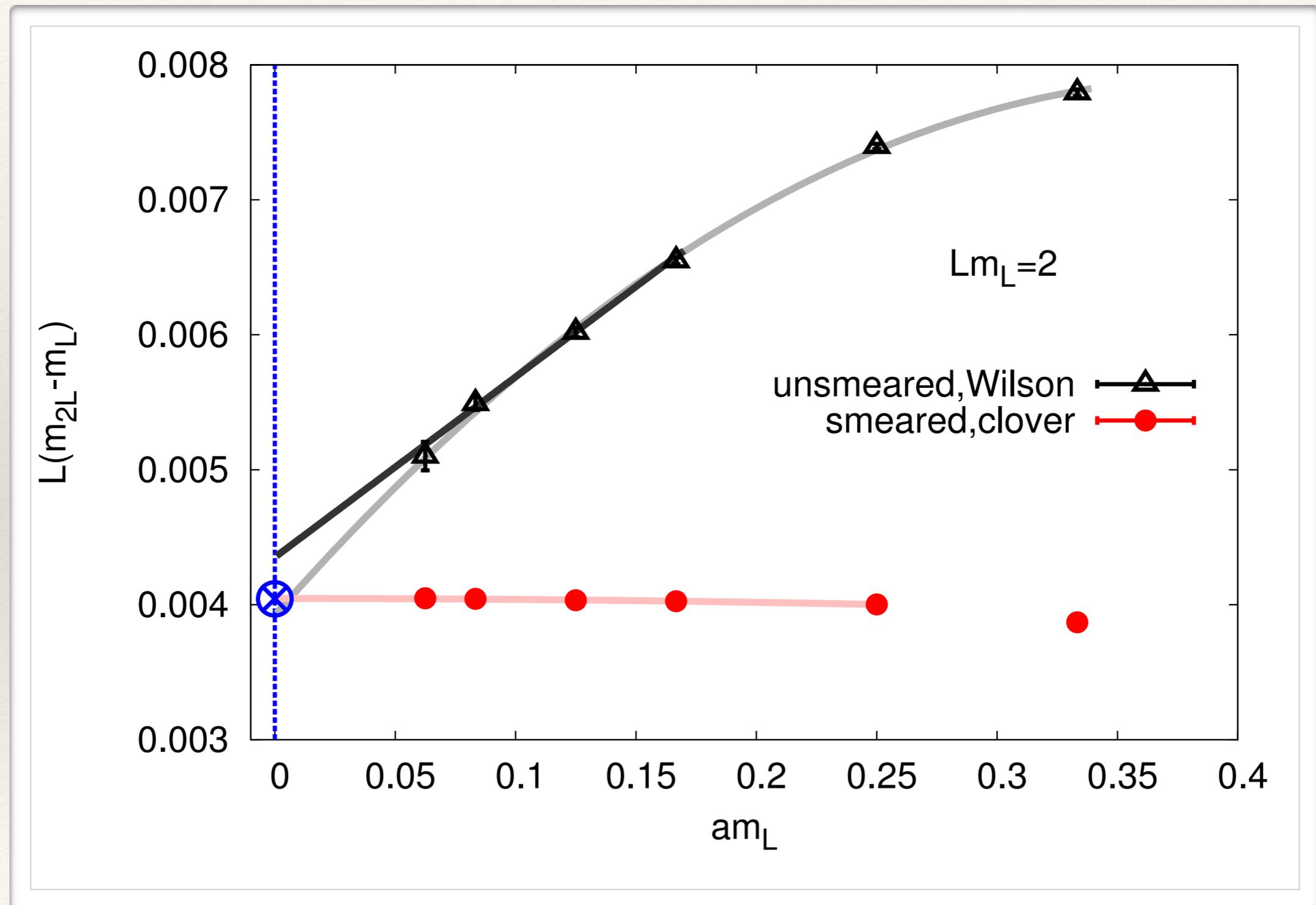
- ❖ Unquenched computations of the light quark masses and Dashen's theorem corrections
- ❖ QCD+QED decay constants are gauge variant and IR divergent. How to deal with that?
C.T. Sachrajda plenary talk: 28/06 — 10:30
- ❖ Compute corrections to matrix elements
($K_{\ell 3}, K \rightarrow \pi\pi, \dots$)
- ❖ QCD+QED to compute hadronic corrections to anomalous magnetic moments.



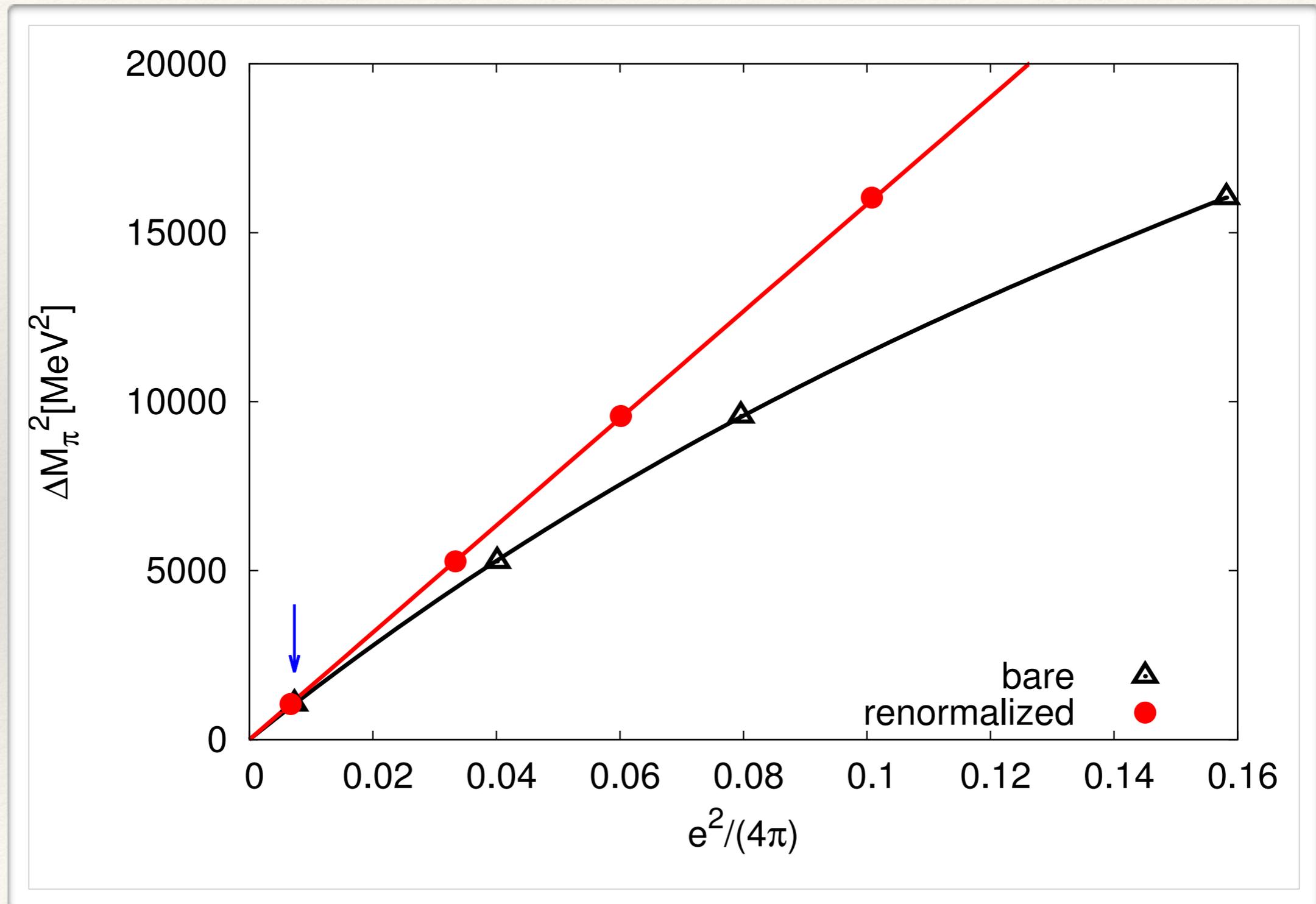
Thank you!

Backup

[BMWc, 2014]: QED simulations



[BMWc, 2014]: charge renormalisation



[BMWc, 2014]: charm discretisation effects

